

# Quantifying simultaneous fluxes of ozone, carbon dioxide and water vapor above a subalpine forest ecosystem

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*Eddy covariance measurements from over a spruce-fir forest are presented and suggestions for improved methods are discussed.*

## Abstract

Assessing the long-term exchange of trace gases and energy between terrestrial ecosystems and the atmosphere is an important priority of the current climate change research. In this regard, it is particularly significant to provide valid data on simultaneous fluxes of carbon, water vapor and pollutants over representative ecosystems. Eddy covariance measurements and model analyses of such combined fluxes over a subalpine coniferous forest in southern Wyoming (USA) are presented. While the exchange of water vapor and ozone are successfully measured by the eddy covariance system, fluxes of carbon dioxide (CO<sub>2</sub>) are uncertain. This is established by comparing measured fluxes with simulations produced by a detailed biophysical model (FORFLUX). The bias in CO<sub>2</sub> flux measurements is partially attributed to below-canopy advection caused by a complex terrain. We emphasize the difficulty of obtaining continuous long-term flux data in mountainous areas by direct measurements. Instrumental records are combined with simulation models as a feasible approach to assess seasonal and annual ecosystem exchange of carbon, water and ozone in alpine environments. The viability of this approach is demonstrated by: (1) showing the ability of the FORFLUX model to predict observed fluxes over a 9-day period in the summer of 1996; and (2) applying the model to estimate seasonal dynamics and annual totals of ozone deposition and carbon, and water vapor exchange at our study site. Estimated fluxes above this subalpine ecosystem in 1996 are: 195 g C m<sup>-2</sup> year<sup>-1</sup> net ecosystem production, 277 g C m<sup>-2</sup> year<sup>-1</sup> net primary production, 535 mm year<sup>-1</sup> total evapotranspiration, 174 mm year<sup>-1</sup> canopy transpiration, 2.9 g m<sup>-2</sup> year<sup>-1</sup> total ozone deposition, and 1.72 g O<sub>3</sub> m<sup>-2</sup> year<sup>-1</sup> plant ozone uptake via leaf stomata. Given the large portion of non-stomatal ozone uptake (i.e. 41% of the total annual flux) predicted for this site, we suggest that future research of pollution–vegetation interactions should relate plant response to actively assimilated ozone by foliage rather than to total deposition. In this regard, we propose the Physiological Ozone Uptake Per Unit of Leaf Area (POUPULA) as a practical index for quantifying vegetation vulnerability to ozone damage. We estimate POUPULA to be 0.614 g O<sub>3</sub> m<sup>-2</sup> leaf area year<sup>-1</sup> at our subalpine site in 1996. Published by Elsevier Science Ltd.

**Keywords:** Ozone deposition; Plant ozone uptake; Carbon dioxide; Flux; Eddy covariance; Latent heat; Sensible heat; Ecosystem model; Ameriflux; Subalpine ecosystem

## 1. Introduction

Forest ecosystems are known to play an important role in carbon sequestration and the uptake and destruction of tropospheric ozone (O<sub>3</sub>). Yet, this role and the budgets of tropospheric carbon dioxide (CO<sub>2</sub>)

and O<sub>3</sub> are uncertain in remote forested areas (Chameides and Lodge, 1992; Wofsy and Hollinger, 1998). With the broad governmental agency support of the new ‘Ameriflux’ program designed to continuously measure ‘net CO<sub>2</sub> flux’ and energy exchange in a variety of environments (Wofsy and Hollinger, 1998), it has become important to choose sites which represent key ecosystem types (such as subalpine areas) and provide valid data. Measurements of ecosystem fluxes at these sites will utilize the eddy covariance technique. Theoretically, this micrometeorological method is only feasible

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over flat, uniform landscapes due to the stationarity and homogeneity assumptions necessary to simplify the turbulent continuity equation for mass. In practice, however, the eddy covariance technique has also been found to produce reliable results in rough terrain (Fitzjarrald and Moore, 1992) when vector rotation is applied (McMillen, 1988).

In 1987, the USDA Forest Service established the Glacier Lakes Ecosystem Experiment Site (GLEES) in southern Wyoming to conduct subalpine ecological research (Musselman, 1994). An Ameriflux installation is currently planned for this study location (Wofsy and Hollinger, 1998). Over the past 10 years, the Forest Service has collected a wealth of information on vegetation and the physical environment at GLEES including data on the ecosystem–atmosphere exchange of mass and energy. Zeller and Hehn (1995, 1996) successfully measured deposition of  $O_3$  and energy fluxes at GLEES in the summer and winter of 1992 and in the winter of 1994. Data on fluxes of  $CO_2$  were obtained at this site in August 1996.

The main objective of this paper is to provide a quantitative assessment of the 1996 seasonal dynamics

and annual totals of key ecosystem fluxes at a subalpine site. We present eddy covariance measurements and model analyses of the concurrent  $CO_2$ ,  $O_3$ , and water vapor fluxes over a high altitude forest at GLEES. We discuss problems related to interpretation of flux measurements in complex terrain using data from a 9-day period in the summer of 1996. Given measurement uncertainties, we explore the possibility of complementing instrumental data records with model simulations using FORFLUX (Nikolov, 1997) to better understand the long-term ecosystem–atmosphere exchange of mass and energy in a subalpine environment.

## 2. Experimental methods

### 2.1. Site description

GLEES is located at 3186 m elevation in the Snowy Range of the Medicine Bow National Forest in southern Wyoming, USA ( $41^{\circ}22'N$ ,  $106^{\circ}14.5'W$ ). The research complex is described in detail by Musselman (1994). The study area is covered by an open forest dominated

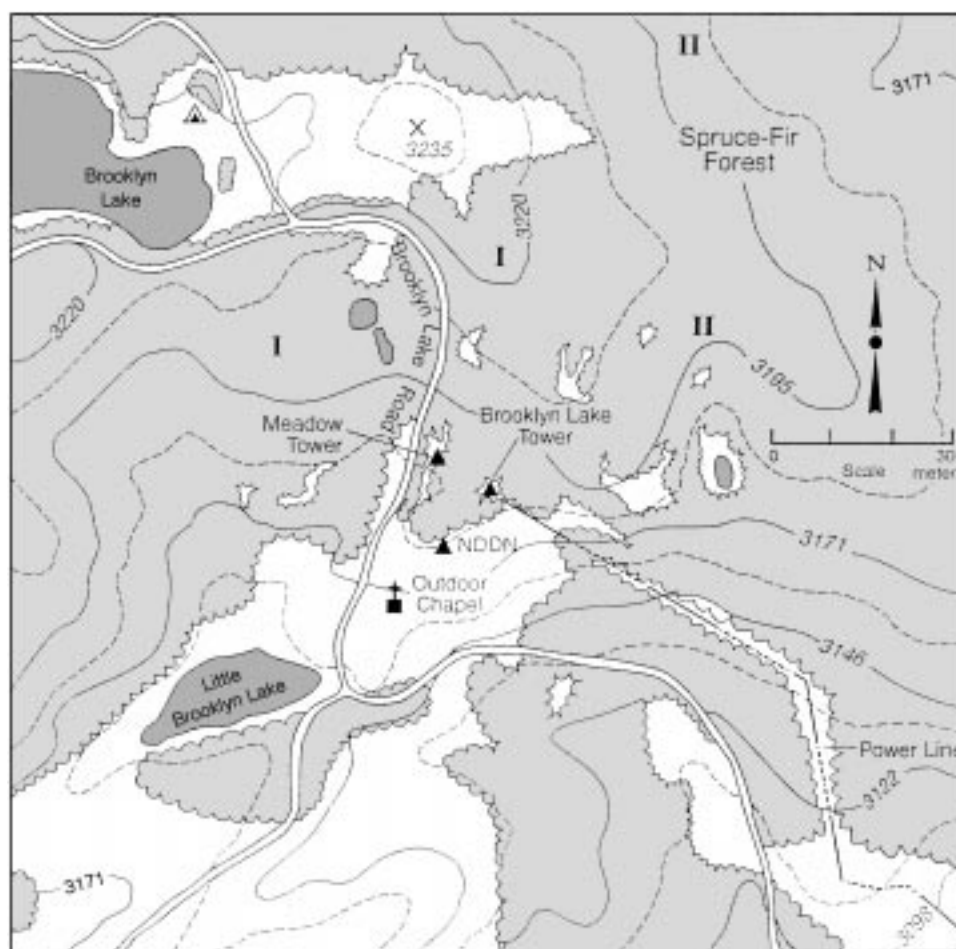


Fig. 1. Site map of Brooklyn Lake tower and surroundings.

by Engelmann spruce (*Picea engelmannii*) and subalpine fir (*Abies lasiocarpa*). The meteorological and flux data were collected from several heights on a 29-m tower (hereafter called the Brooklyn tower) situated in a 30-m wide forest opening (Fig. 1). Wind and CO<sub>2</sub> data were also collected from 1-m height at a 'wet' meadow site located 100 m west-northwest of the Brooklyn tower. This small tower is positioned at the east edge of a meadow and is separated from the Brooklyn tower by a uniform flat forested area. Measurements in the meadow are used to characterize the near-surface CO<sub>2</sub> concentration and wind field around the Brooklyn tower. The average height of the forest canopy at the Brooklyn tower is 17 m. For a neutrally stable atmosphere the calculated displacement height ( $d$ ) is 11.7 m and the roughness length ( $z_o$ ) is 1.7 m (Zeller and Hehn, 1996). For modeling purposes, we have estimated the leaf area index (LAI) of the surrounding forest to be 2.8 m<sup>2</sup> m<sup>-2</sup>.

Within 1 km of the tower, the terrain has a slope of +2.5% from west to east and -9.7% from north to

south. In the direction of the predominant wind (i.e. west-northwest), the landscape is forested, and has a constant slope of 2.2% for about 1 km. The predominant wind direction above the canopy is 210° (south-southwest)–330° (northwest). Fig. 2 illustrates the effect of terrain on the half-hour wind vector ratio ( $w/u$ ) at 23 and 1 m above the ground as a function of wind direction. It can be seen that at 23 m, the wind first flows uphill (positive  $w$ ) from south-southwest through west-northwest, and then abruptly flows downhill as the wind direction swings farther from north. The topographic contour lines in Fig. 1 show a small ridge just east of the tower that causes the vertical flow change. The wind vector ratios measured at 1 m height on the meadow tower are predominantly negative. They are also plotted on Fig. 2 and discussed below.

This site is relatively complex from a standpoint of micrometeorological measurements. CO<sub>2</sub> has not been routinely measured at GLEES but O<sub>3</sub> concentrations have been documented by Wooldridge et al. (1997) and O<sub>3</sub> fluxes by Zeller and Hehn (1996).

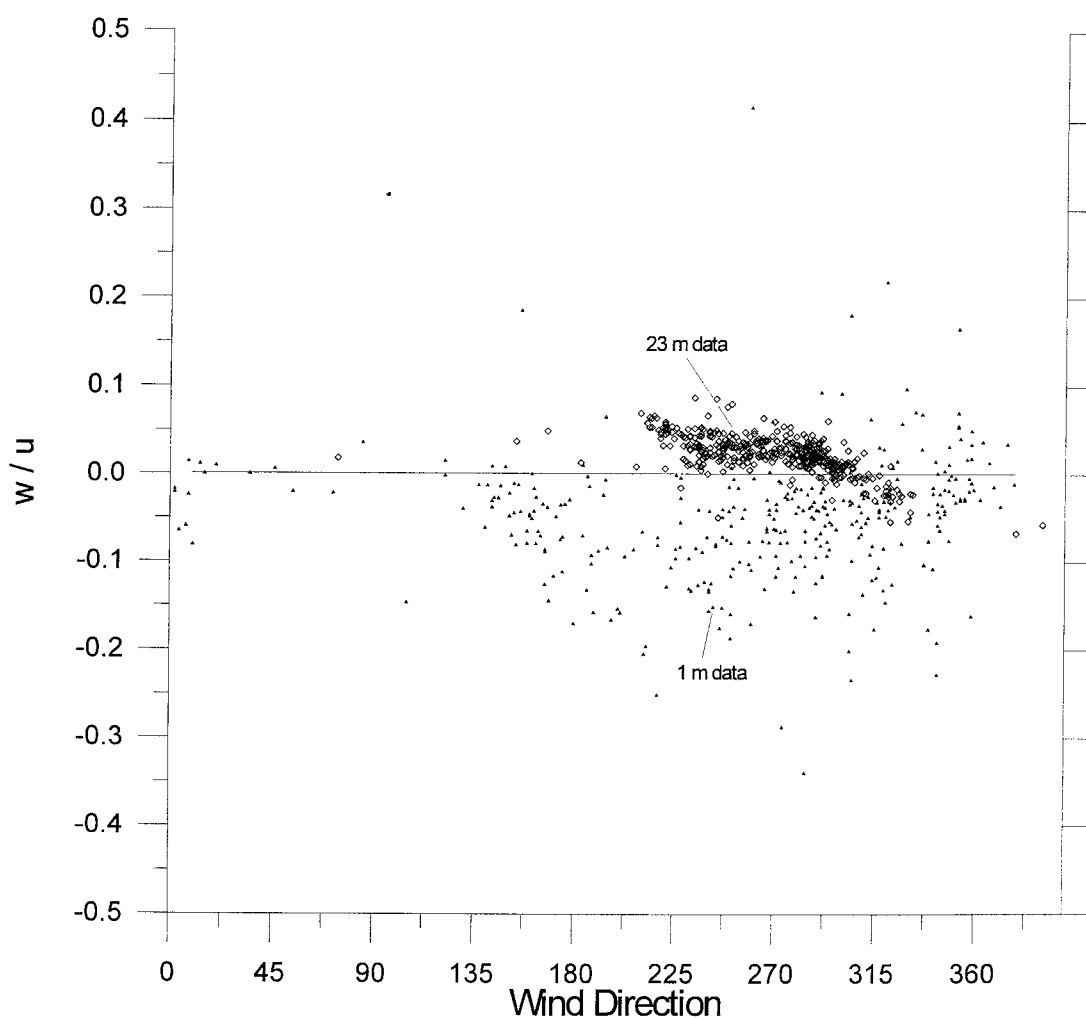


Fig. 2. Plot of half-hour wind vector magnitude ratio ( $\overline{w/u}$ ) as a function of wind direction for the period 30 July–8 August, 1996. Open diamonds ( $\diamond$ ) are 23 m data and filled triangles ( $\blacktriangle$ ) are 1 m data. (A value of  $w/u=0.1$  is equivalent to 5.7° vertical angle.).

Table 1  
List of eddy covariance instrumentation at Brooklyn Lake meadow site

Instrument or model	Height(s)	Type	Parameter	
TECO49 O <sub>3</sub> Analyzer	29	UV adsorption	$c$	O <sub>3</sub> concentration
Infrared gas analyzer	23 1	Open-path, infrared absorption	$c, c'$ $q, q'$	CO <sub>2</sub> and H <sub>2</sub> O concentration and fluctuation
CAAM1 O <sub>3</sub> Analyzer	29	Chemiluminescence (fast response)	$c'$	O <sub>3</sub> fluctuation
RM Young Gill anemometer	23	Propeller	$w, w'$	Vertical wind
	29	Anemometer	$u, u'$	Horizontal wind
			$v, v'$	Horizontal wind
ATI, Inc. 15-cm anemometer	1	Sonic anemometer	$w, w'$	Vertical wind
			$u, u'$	Horizontal wind
			$v, v'$	Horizontal wind
Campbell Sc., Inc. krypton hygrometer	29	Light adsorption	$q, q'$	H <sub>2</sub> O

## 2.2. Instrumentation and flux measurements

Micrometeorological and flux data were collected from 30 July to 8 August, 1996. This data represent 10 out of the 13 ‘core’ measurement parameters and three out of the seven ‘desired’ parameters recommended by the Ameriflux science team (Table 2 in Wofsy and Hollinger, 1998).

The eddy covariance systems (Zeller et al., 1989; Massman et al., 1990) were deployed at 23 (6 m above canopy) and 29 m (12 m above canopy) on the Brooklyn tower. They measured fluxes of O<sub>3</sub>, CO<sub>2</sub>, momentum, sensible and latent heat as well as concentrations of O<sub>3</sub> and CO<sub>2</sub>, temperature, wind speed and wind direction. Standard meteorological sensors were permanently mounted at 10 and 29 m for routine weather measurements (Musselman, 1994). O<sub>3</sub> flux was measured at 29 m employing a Gill<sup>1</sup> UVW anemometer and a Chemiluminescence Ambient Air Monitor (CAAM1) (Ray et al., 1986). CAAM1 was continuously calibrated using a TECO49 commercial UV adsorption instrument. For O<sub>3</sub> measurements, ambient air was sampled from 29 m height through a 31-m long (1.6-cm inside diameter) Teflon tube. The lag time of the intake system ( $t_1$ ) was typically 2.4 s. A second TECO49 instrument was operated at 3 m height near the base of the tower to measure below-canopy O<sub>3</sub> concentration. Fluxes of CO<sub>2</sub> and water vapor were measured at 23 m on the Brooklyn tower using an open-path, infrared absorption gas analyzer (Auble and Meyers, 1992). CO<sub>2</sub> measurements were also made at 1 m height on the meadow tower using a second open-path CO<sub>2</sub> analyzer. The slow responding CO<sub>2</sub>–H<sub>2</sub>O analyzer (LICOR 6262), used for continuous calibration of the CO<sub>2</sub> instrument deployed

at 23 m, malfunctioned just prior to the 9-day period presented here. The Gill UVW anemometers have been used on the Brooklyn tower instead of sonic anemometers for their greater durability in a harsh alpine environment and were used in this study because two sonics were not available. The measurements taken at the meadow at 1 m height employed a 15-cm sonic anemometer, and a second open-path, infrared absorption CO<sub>2</sub> gas analyzer. Table 1 summarizes the flux instrumentation used in this study.

The 15-Hz eddy deviations of O<sub>3</sub> and CO<sub>2</sub> concentrations,  $c' = c - \bar{c}$  (where  $\bar{c}$  is a 200 s recursive filter average concentration) were multiplied by the turbulent wind components  $u'$ ,  $v'$  and  $w'$ , and then averaged over a halfhour sampling period (McMillen, 1988) to obtain the vertical flux ( $F_c$ ) ( $\mu\text{g m}^{-2} \text{s}^{-1}$ ) after a vector coordinate rotation for the  $\bar{w}=0$  streamline, i.e.:

$$F_c = \overline{w'(t)c'(t - t_1)} \quad (1)$$

Here, negative  $F_c$  indicates downward flux. Fluxes of momentum, sensible and latent heat were obtained likewise. The coordinate rotation does not affect the scalar flux sign and only slightly influences measured flux magnitudes. Fluxes with associated streamlines within  $\pm 3^\circ$  (at 23 m) and  $\pm 8^\circ$  (at 29 m) of the horizontal account for 98% of the data presented here (note that  $w/u=0.1$  is equivalent to  $5.7^\circ$  in Fig. 2). Gill UVW anemometer data were corrected in real time for the inherent cosine response (Massman and Zeller, 1988).

Businger’s list of measurement concerns (Businger, 1986) was adopted for evaluation and editing of the eddy covariance data. The O<sub>3</sub> flux corrections for sensible and latent heat effects implemented Eq. (12) by Leuning and Moncrieff (1990). The corrections of CO<sub>2</sub> flux for an open-path system were made using Eq. (5) by Leuning and Moncrieff (1990). Hereafter, we refer to

<sup>1</sup> The use of trade and company names is for the benefit of reader; such use does not constitute an official endorsement or approval of any service or product by the USDA to the exclusion of others that may be suitable.

the latter correction as LM. We discuss the effect of these corrections below in addition to an empirical ‘advection’ correction based on Lee (1998) and Baldocchi (personal communication). The latent heat flux obtained with a krypton hygrometer at 29 m height was corrected for the oxygen influence (Tanner and Greene, 1989). Data below the scale height of 23 m (i.e.  $h_t < |w'\bar{c}'/(\partial c/\partial t)|$ ) for CO<sub>2</sub> and below 29 m for O<sub>3</sub> were culled. A digital butterworth filter was applied in real time to account for aliasing. Corrections for instrument design and separation can theoretically increase flux magnitudes by 20–80% in uniform flat terrain (Zeller et al., 1989; Massman et al., 1990) when using a Gill anemometer with faster responding instruments (e.g. the CO<sub>2</sub> and H<sub>2</sub>O sensors). These corrections tend to increase the random noise in reported fluxes, but do not affect flux direction and are much less important in the above-canopy roughness layer where transporting eddies are relatively large. Hence, the flux data presented here were not corrected for such effects.

Neutral stability statistics (Panofsky and Dutton, 1984; Schuttlesworth, 1989) during this study for wind, temperature, CO<sub>2</sub> and O<sub>3</sub> are fairly consistent and are equivalent to results from other field studies. For example,  $\sigma_w/u^*$  was consistently  $1.2 \pm 0.1$  at 29 m and  $0.95 \pm 0.1$  at 23 m compared to 1.3 for flat terrain. Dimensionless wind shear,  $\phi_m = \overline{kz/u^*}(\partial u/\partial z)$ , was consistently  $1.3 \pm 0.3$  at 29 m and  $1.1 \pm 0.3$  at 23 m where  $u^* = \sqrt{-w'u'}$  was calculated from the momentum flux.

### 2.3. Overview of the FORFLUX model

Nikolov (1997) developed a biophysical model (FORFLUX) to study the seasonal exchange of CO<sub>2</sub>, O<sub>3</sub>, and energy between terrestrial ecosystems and the atmosphere. The model mechanistically couples major ecosystem processes controlling the flows of carbon and water implementing recent concepts in plant eco-physiology and soil hydrology. FORFLUX consists of four interconnected modules: a leaf photosynthesis model, a canopy flux model, a soil heat-, water- and CO<sub>2</sub>-transport model, and a snow pack model.

Leaf-level net photosynthesis and water-vapor flux are computed by the LEAFC3 submodel (Nikolov et al., 1995). This submodel simulates coupled interactions between mesophyll biochemical reactions, stomatal conductance and the leaf energy balance. LEAFC3 predicts four state variables, which are used as input to the FORFLUX canopy module, i.e. net photosynthesis rate, stomatal conductance to water vapor, leaf temperature and leaf boundary-layer conductance. These variables are obtained from the solution to a system of seven simultaneous equations.

The FORFLUX canopy module scales LEAFC3 output to a canopy level. This is achieved by numerical integration of the LEAFC3 model over the canopy LAI.

The integration takes into account: (1) the short- and long-wave radiative transfer inside the canopy; (2) variation of foliage photosynthetic capacity with canopy depth; (3) wind speed attenuation throughout the canopy; and (4) rainfall interception by foliage elements. The canopy submodel also calculates above-ground woody respiration and seasonal phenological changes in stand LAI. The radiative transfer module implements a four-stream analytical solution to the multiple scattering equation and is based on work by Ross (1981), Sellers (1985) and Camillo (1987). The algorithm accounts for spectral changes of the incident solar flux inside the canopy caused by the selective radiation absorption of green leaves. The model uses leaf clumping and angular distribution of foliage elements to estimate penetration and absorption of the direct and diffuse radiation in the visible (380–700 nm), near-infrared (700–1500 nm) and thermal (>1500 nm) bands. Leaf photosynthetic capacity is assumed to decline with canopy depth according to data and models by Hirose and Werger (1987) and Hirose et al. (1988). The attenuation of horizontal wind speed inside the canopy is simulated using the phenomenological model of Albin (1981) as modified by Massman (1987). In FORFLUX, wind speed controls the energy balance of leaves by modifying the leaf-boundary layer conductance and air humidity at the leaf surface. Rainfall interception is computed as a function of canopy water storage capacity, rainfall intensity and current evaporative demand employing an analytical interception model by Massman (1983). Net photosynthesis of wet foliage is estimated by the LEAFC3 model assuming that evaporation completely suppresses transpiration, and wetness does not adversely affect CO<sub>2</sub> assimilation capacity of leaves.

The soil module simulates the dynamics of moisture and temperature within the soil profile, soil-surface evaporation and CO<sub>2</sub> efflux from soil due to root and microbial respiration. Vertical transport of water and heat in soil is modeled using principles of diffusion theory (Hillel, 1982; Campbell, 1985; Feddes et al., 1988). The effect of soil heterogeneity on field-scale fluxes is simulated employing the Bresler–Dagan concept (Bresler and Dagan, 1983; Dagan and Bresler, 1983; Dagan, 1993). The latter can be summarized as follows. A spatially variable field is approximated by an array of non-interacting homogeneous soil columns, each one of which is characterized by a unique set of hydraulic parameter values. These values are estimated from observed probability distributions of soil properties across the horizontal plane. The hypothesis is that one can obtain an accurate estimate of the mean scalar fluxes at the field scale by computing scalar transports in numerous columns (using a 1-dimensional model) and then averaging the output horizontally. The theoretical justification behind such an approach is that although soil

properties vary in a 3-dimensional space, the length scale of variation is much larger in the horizontal domain than it is in the vertical one. Thus, vertical soil variability would only have a small effect on the average field flow. Soil surface evaporation is computed from an explicit analytical solution of the surface energy balance equation, which accounts for atmospheric stability effects in the surface boundary layer. The CO<sub>2</sub> efflux from soil is estimated using a bulk-parameterization approach, which considers the effect of soil moisture, temperature and clay content on CO<sub>2</sub> production. The temperature factor is computed using the exponential function of Kirschbaum (1995) and soil respiration data summarized by Peterjohn et al. (1994). The effect of volumetric moisture on soil CO<sub>2</sub> release is described using experimental data and models by Schlentner and Van Cleve (1984) and Grant and Rochette (1994).

The snowpack module accounts for the dynamics of snow accumulation and melt on the ground. Incoming precipitation is partitioned into snow- and rainfall based on ambient temperature (Wigmosta et al., 1994). The model does not consider snow interception by the canopy and, hence, all snowfall is added to the ground storage. Snowpack is depleted through sublimation and melting. Both processes are assumed to occur at the snow–atmosphere interface and to be a function of the snowpack energy balance (Male and Gray, 1981). Radiation absorption by the snowpack is computed from incident visible, near-infrared and thermal irradiances and corresponding snow reflectances. Snow albedos in the visible and near-infrared wavelength are estimated using the physical model of Marshall (1989), which accounts for effects of snow grain radius, impurity content, snow depth, solar zenith angle and cloud cover. Snow depth and soot concentration primarily influence visible albedo while grain size and zenith angle mainly affect near-infrared snow reflectance.

All biophysical interactions in the FORFLUX model are computed hourly while model projections are produce on hourly or daily time step. FORFLUX requires for input hourly values of ambient temperature, relative humidity, incident short-wave radiation, precipitation and above-canopy wind speed. Species physiology is characterized in the model by 20 parameters. Tables A1 and A2 in the Appendix list the species and site-specific input parameters and their values used in this study. Model output includes temporal courses of net ecosystem carbon and water fluxes and their components, such as vegetation net primary production (NPP), canopy photosynthesis and stomatal conductance, woody respiration, soil CO<sub>2</sub> efflux, canopy transpiration and rainfall interception, soil evaporation, snow melt and sublimation, surface runoff and subsoil drainage.

Compared to other biophysical models such as BIOMASS (McMurtie et al., 1990, 1992), FOREST-BGC

(Running and Coughlan, 1988; Running and Gower, 1991), PnET (Aber and Federer, 1992), and TCX (Bonnan, 1991a,b), the FORFLUX model is notable for the following innovations:

1. Canopy fluxes are estimated by a vertical integration of the LEAFC3 model (Nikolov et al., 1995) which mechanistically couples photosynthesis, stomatal conductance and energy balance at the leaf level. This provides a more accurate simulation of the gas and energy exchange at the canopy level compared to other models, which do not incorporate feedbacks between latent heat flux and CO<sub>2</sub> uptake.
2. The transport of water and heat in soil is modeled employing physical principles of diffusion theory (Feddes et al., 1988) and explicitly accounting for effects of soil heterogeneity on field-scale fluxes (Dagan, 1993). This approach is expected to yield predictions of soil moisture dynamics, which are far more robust than those provided by various soil-bucket approximations used in other ecosystem models. To our knowledge, FORFLUX is the only biophysical model which combines a detailed formulation of canopy processes with a physically based stochastic algorithm for simulating heat and water flow in soil.
3. Root–shoot communication is modeled using a new scheme, which incorporates the joined effect of root chemical and hydraulic signaling on leaf stomatal conductance (Tardieu and Davies, 1993). This allows for a realistic simulation of the observed feedback between transpiration rate and stomatal sensitivity to adverse soil conditions such as drought and freezing.

### 2.3.1. Simulation of O<sub>3</sub> deposition

The FORFLUX model predicts total O<sub>3</sub> deposition ( $F_{O_3}$ ) as a sum of three fluxes, i.e. plant O<sub>3</sub> uptake via leaf stomata, O<sub>3</sub> deposition to non-transpiring plant surfaces, and O<sub>3</sub> flux absorbed by the ground (i.e. soil or snow surface), i.e.:

$$F_{O_3} = \frac{\Delta C_{O_3}}{r_a + 1/g_{cs}} + \frac{\Delta C_{O_3}}{r_a + 1/g_{ct}} + \frac{\Delta C_{O_3}}{r_a + r_u + r_{bs} + r_{is}}. \quad (2)$$

In this equation,  $\Delta C_{O_3}$  is the O<sub>3</sub> concentration gradient (ppbv) between the active surface and a reference height above the canopy (assuming a zero O<sub>3</sub> concentration at the surface;  $\Delta C_{O_3}$  equals the measured ambient O<sub>3</sub> concentration with a negative sign);  $g_{cs}$  is the foliage canopy conductance to O<sub>3</sub> (m s<sup>−1</sup>) (i.e. the sum in parallel of stomatal and leaf boundary-layer conductances);  $g_{ct}$  is the plant surfacial (or cuticular) conductance to O<sub>3</sub> deposition (ms<sup>−1</sup>);  $r_a$  and  $r_u$  are the

above- and within-canopy aerodynamic resistances, respectively;  $r_{bs}$  is the soil boundary-layer resistance; and  $r_{is}$  is the soil intrinsic resistance (all resistances have units of  $s\ m^{-1}$ ). The canopy  $O_3$  conductance ( $g_{cs}$ ) is the most important one of all conductances and resistances within the soil–plant–atmosphere continuum. It is estimated by integrating individual leaf  $O_3$  conductances ( $g_{lt}$ ) over the canopy LAI ( $L_t\ m^2\ m^{-2}$ ), i.e.:

$$g_{cs} = \int_0^L g_{lt} dL, \quad (3)$$

where  $g_{lt}$  is calculated from the leaf stomatal ( $g_{sv}$ ) and boundary-layer ( $g_{bv}$ ) conductance to water-vapor exchange using the formula:

$$g_{lt} = \frac{g_{sv}g_{bv}}{1.315g_{sv} + 1.508g_{bv}}. \quad (4)$$

In Eq. (4), 1.508 and 1.315 are ratios of water vapor and  $O_3$  diffusivities in still air and in the leaf-boundary layer, respectively. These are calculated from data presented in Table 2 by Massman (1998). Vapor conductances, on

the other hand, are predicted by the coupled leaf photosynthesis–transpiration module (LEAFC3) of the FORFLUX model.

Canopy aerodynamic resistance ( $r_a$ ) is estimated using an empirical equation derived from the measured momentum flux at GLEES:

$$r_a = \frac{19}{U^{1.35}}, \quad (5)$$

where  $U$  is the measured horizontal wind speed ( $m\ s^{-1}$ ) above the canopy. The plant surfacial conductance is estimated as  $g_{ct} = 1.4 \times 10^{-5}\ m\ s^{-1}$  based on data by Kerstiens and Lenzian (1989) and Massman and Grantz (1995). The within-canopy aerodynamic resistance is assumed to be eight times greater than  $r_a$ , i.e.:

$$r_a = 8r_{a_i}. \quad (6)$$

The surface boundary-layer resistance to  $O_3$  is computed from the ground boundary-layer conductance to water vapor ( $g_{bs,v}\ m\ s^{-1}$ ) using the formula:

$$r_{bs} = \frac{1.37}{g_{bs,v}}, \quad (7)$$

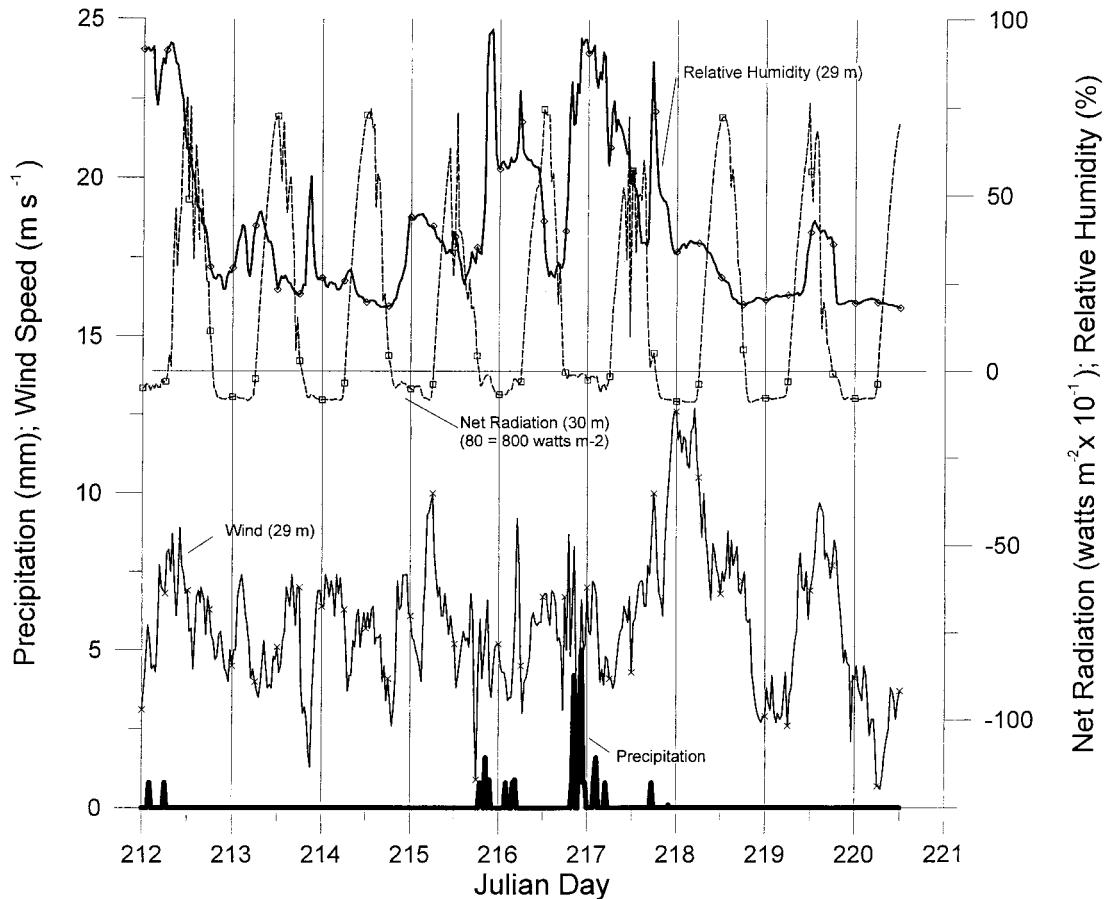


Fig. 3. Upper: half-hour average relative humidity (solid line, open diamonds); half-hour net radiation (dashed line, open squares). Lower: half-hour wind speed at 29 m (solid line, ex's); precipitation (thick solid line) for the period 30 July–8 August, 1996.

where  $g_{bs,v}$  is produced by the ground evaporation module of the FORFLUX model. This module solves exactly the surface energy balance equation while accounting for atmospheric stability effects on the boundary-layer conductance. The constant 1.37 is a ratio of molecular diffusivities of  $O_3$  and water vapor assumed in the surface boundary layer.

The intrinsic soil resistance is considered to be constant in time. Based on information on soil resistances to  $O_3$  deposition gathered by Turner et al. (1974) and van Pul (1992), it was assumed  $r_{is} = 410 \text{ s m}^{-1}$  for bare soil. In the presence of a snow pack,  $r_{is}$  is set to  $6000 \text{ s m}^{-1}$  as suggested by measurements of Zeller and Hehn (1995).

### 3. Results and discussion

#### 3.1. Surface Meteorology

Fig. 3 depicts the 9-day time series of measured half-hour values of net radiation, relative humidity, wind

speed and precipitation. Air temperature is graphed on the bottom of Fig. 4. Four short precipitation events occurred around midnight on Julian days 212, 216, 217 and 218. Rainfall interrupts the fast response open-path sensors. Therefore,  $CO_2$  fluxes during these periods are automatically eliminated. The use of Gill anemometers and the air intake system for  $O_3$  allows continuous sampling of  $O_3$  fluxes during light rainfall. We assume that light rainfall in the turbulent canopy boundary layer has a negligible effect on measured fluxes and, thus, do not correct for it. Since sensible and latent heat fluxes are intermittently affected by precipitation, hours having abnormally high or low variances are also routinely culled (i.e. sensible:  $> 1.2^\circ C^2$  and  $< 0.0001^\circ C^2$ ; latent:  $> 0.4 \text{ g}^2 \text{ m}^{-6}$  and  $< 0.0001 \text{ g}^2 \text{ m}^{-6}$ ).

#### 3.2. $O_3$ and $CO_2$ concentrations

Temporal variations of  $O_3$  and  $CO_2$  concentrations are presented in Fig. 4. There is some diurnal correlation between temperature and  $O_3$  concentration

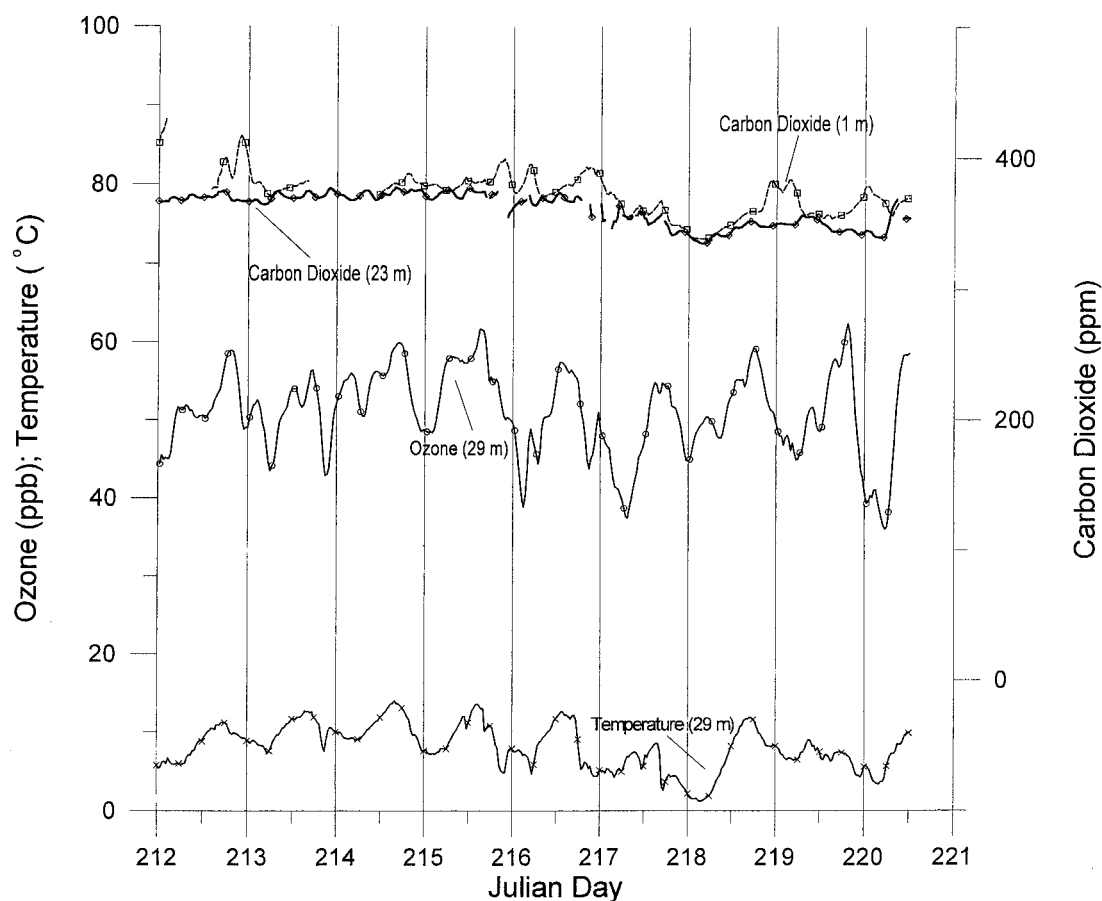


Fig. 4. Half-hour average  $CO_2$  concentration (23 m solid line, open diamonds; 1 m  $O_3$  concentration (dashed),  $O_3$  concentration (solid line, open circle) and temperature (solid line ex's) for the period 30 July–8 August, 1996.



typically observed with  $O_3$  monitoring.  $CO_2$  concentrations are less accurate as they are obtained by highly sensitive fast-response sensors. However, the measured values average 377 ppm over the 9-day period. Half-hour variations of concentration at 1 and 23 m height are noisier than typically measured with slower-response instruments.

### 3.3. Measured vertical gradients

The measured vertical gradients of  $O_3$ ,  $CO_2$  and temperature are portrayed in Fig. 5. Positive values indicate increase of concentration with height. Thus, a positive gradient is associated with a negative (downward) flux, and a negative gradient implies a positive (upward) flux. In accordance with expectations, temperature gradients are greater near the surface as seen from the comparison of records from 10 and 3 m, and between 29 and 10 m height. The effect of enhanced night-time turbulent mixing on gradients of temperature and  $O_3$  concentrations can be seen on days 215 and 218 when high wind speeds at night cause measured gradients to approach zero. To some extent this is also true for  $CO_2$

concentrations. The  $CO_2$  gradient is calculated from the open-path sensors, which are separated vertically by 22 m and horizontally by 100 m.

### 3.4. Energy fluxes and energy balance

The diurnal dynamics of momentum, sensible and latent heat fluxes are presented in Fig. 6 along with the soil heat flux. The 24-h average residuals of the daily energy balance ( $R_n - G - H - LE$ ) ending at midnight on each day are plotted on Fig. 6. These values are positive through JD 215, missing JD 215 through 217 due to night-time precipitation, and are positive but lower in value after JD 218. Daily residuals closely follow the temperature record for the sampling period (Fig. 4). This might indicate that the ecosystem is storing and releasing energy on a scale of several days. The effect of a possible energy storage and advection (Lee, 1998) can also be seen on the energy balance plot for the 9-day period (Fig. 7). Note that the TERM ( $R_n - G$ ) is larger (positive values) than ( $H + LE$ ) during daytime and smaller (negative values) at night.

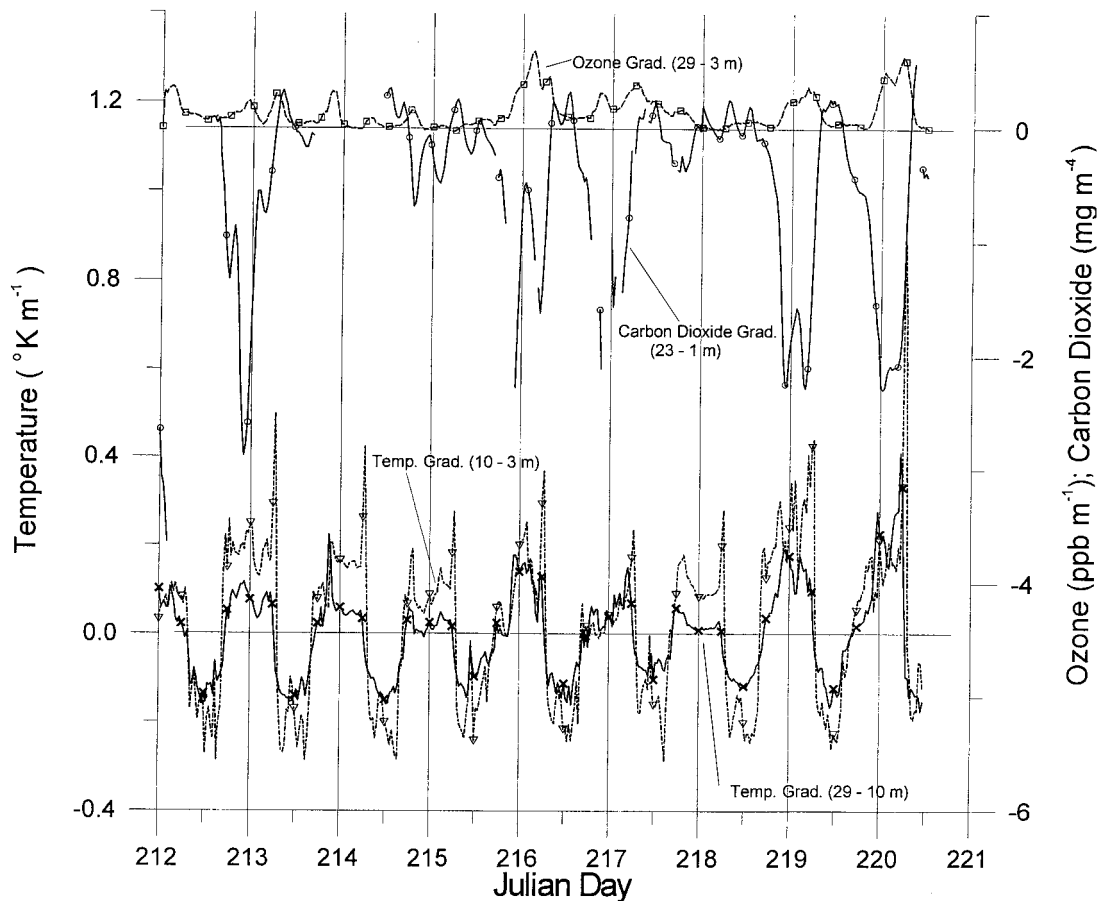


Fig. 5. Half-hour average 29–3 m  $O_3$  gradient; 23–1 m  $CO_2$  gradient; and 29–10 m, and 10–3 m temperature gradients for the period 30 July–8 August, 1996.

### 3.5. $O_3$ and $CO_2$ fluxes

The eddy covariance measurements shown in Fig. 8 are consistent with respective vertical gradients (Fig. 5). The local time rate of  $CO_2$  concentration change,  $\int_0^h \frac{\partial c}{\partial t} \delta z$ , in the canopy layer was estimated from measured data as  $((\Delta \bar{c} / \Delta t)_{h2} - (\Delta \bar{c} / \Delta t)_{h1}) \Delta z$ , where  $h=1$  and 23 m for  $CO_2$ , and 3 and 29 m for  $O_3$ . This estimate, also considered a storage term (Lee, 1998), is plotted for  $CO_2$  and  $O_3$  in Fig. 8. When this value was significant compared to measured  $O_3$  eddy flux (e.g.  $>20\%$ ), the stationarity assumption in Eq. (1) was judged to be violated and the associated eddy flux term was culled. However, for the  $CO_2$  flux analysis discussed below, we kept track of this storage term and did not cull the corresponding  $CO_2$  flux unless it was greater than 100%. Unfortunately, the corrected  $CO_2$  flux (in Fig. 8) is not continuous due to either lack of information to complete the corrections or due to the presence of interfering atmospheric conditions. Such data losses are inevitable when using the eddy covariance technique and give an idea of the data gaps one can expect from  $CO_2$  flux measurements in rough terrain.

#### 3.5.1. Carbon dioxide fluxes

The  $CO_2$  fluxes are presented in Fig. 8 as three superimposed data records. The three records illustrate the analyses problems associated with using the eddy covariance technique to measure  $CO_2$  exchange in complex terrain. The thin line represents the calibrated 'raw' flux as measured prior to corrections. The darker dashed line (only discernible during night-time hours) shows the result of applying the LM correction for open-path sensors. This correction is very large. It varies from  $+0.5 \pm 0.1 \text{ mg m}^{-2} \text{ s}^{-1}$  during daytime to  $-0.15 \pm 0.05 \text{ mg m}^{-2} \text{ s}^{-1}$  at night for the sensible heat term, and from  $+0.08 \pm 0.03 \text{ mg m}^{-2} \text{ s}^{-1}$  (daytime) to  $-0.002 \pm 0.002 \text{ mg m}^{-2} \text{ s}^{-1}$  (night-time) for the latent heat term. The sensible heat term in the LM correction accounts for 85% of the adjustment and the latent heat term for the remaining 15%. The  $CO_2$  flux resulting from the LM correction is near zero for most nights and near zero and random during daylight for days 215–219. For the  $CO_2$  fluxes measured by eddy covariance in complex terrain, Lee (1998) suggests an advection correction of the form  $\overline{w_r}(\bar{c}_r - \langle c \rangle)$  and a storage correction to account for the night-time loss of  $CO_2$  below the flux

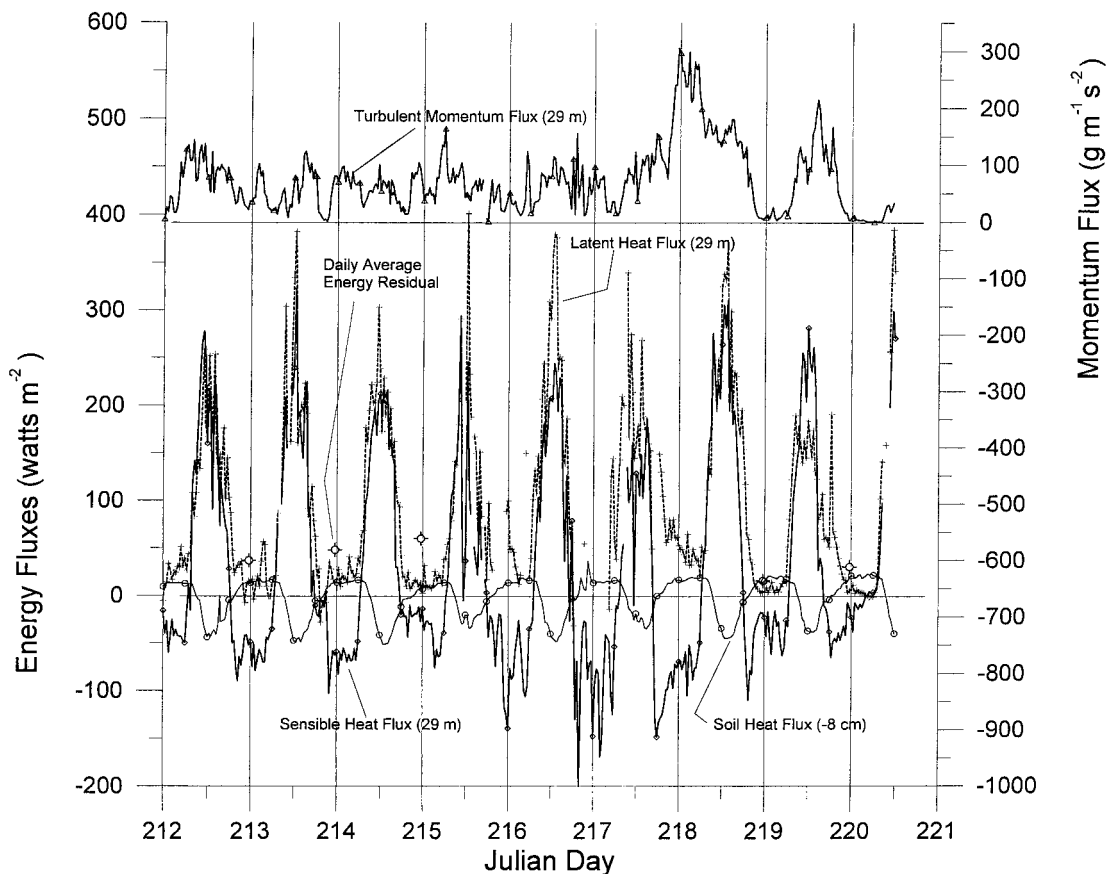


Fig. 6. Half-hour average 29 m turbulent momentum flux (Reynold's stress) (solid line, open triangle); 29 m sensible heat energy flux (solid line, open diamond); 29 m latent heat energy flux (dashed line, plus sign); -8 cm soil heat flux (thin solid line, open circle); and 24 h energy balance average (plus sign superimposed with open circle) for the period 30 July–8 August, 1996.

sampling height. Baldocchi (personal communication) suggests the correction  $\overline{w_a}(\overline{c_r} - \langle c \rangle)$ . Here  $w_a$  is the difference between the measured vertical velocity,  $w_r$ , and a mean vertical velocity ( $w_\theta$ ) which depends on terrain and wind direction (i.e.  $w_a = w_r - w_\theta$ );  $c_r$  is the scalar concentration at reference height; and  $\langle c \rangle$  is the mean scalar concentration in the layer below the sampling height. Based on the observed large  $\text{CO}_2$  gradients (Fig. 5), and the small magnitude of the corrected night-time flux, we conclude that the true  $\text{CO}_2$  flux is not seen at 23 m and that an advection phenomena is occurring at our Brooklyn site. Due to the random fluctuation of the  $\text{CO}_2$  storage term, the latter was not included in the correction as suggested by Lee (1998).

We attempted an advection correction (Lee, 1998; Baldocchi, personal communication). To this end, we used data on vertical wind velocity ( $w_r$ ) measured at 23 m (Fig. 2) to compute a wind direction–terrain dependent vertical velocity ( $w_\theta$ ) to obtain  $w_a$ . Even though the term ( $c_r - \langle c \rangle$ ) is consistently negative at night and positive during daylight,  $w_a$  at 23 m fluctuated widely making

the correction  $\overline{w_a}(\overline{c_r} - \langle c \rangle)$  vary randomly between  $\pm 0.2 \text{ mg m}^{-2} \text{ s}^{-1}$ . We also tried to use vertical wind velocity recorded at 1 m at the meadow site to estimate  $w_a$  but obtained a similarly unsatisfactory fluctuating result. Nevertheless, we judge that for our site, the wind data measured at 1 m better represent the near surface advection of air and  $\text{CO}_2$  compared to those obtained at 23 m. Baldocchi (personal communication) suggested that  $w_\theta$  and  $\overline{w_a}(\overline{c_r} - \langle c \rangle)$  should be established as averages using a long time series and that the correction is best applied as a diurnal function. Without the benefit of a complete  $\text{CO}_2$  profile, but based on the  $\text{CO}_2$  gradient in Fig. 5 and the slower wind speed measurements at 1 m, we hypothesize that  $\text{CO}_2$  is most likely being confined to the lower 1–2 m atmospheric layer. We also chose to use the friction velocity ( $u^*$ ) at 1 m as a parameter to advect  $\text{CO}_2$  in the surface layer since it is an indication of both air movement and turbulence. Hence, we employed  $-u^*(\overline{c_r} - \langle c \rangle)$  as an empirical estimate for the advection. The result from this correction plus the LM correction is shown as a heavy solid line in

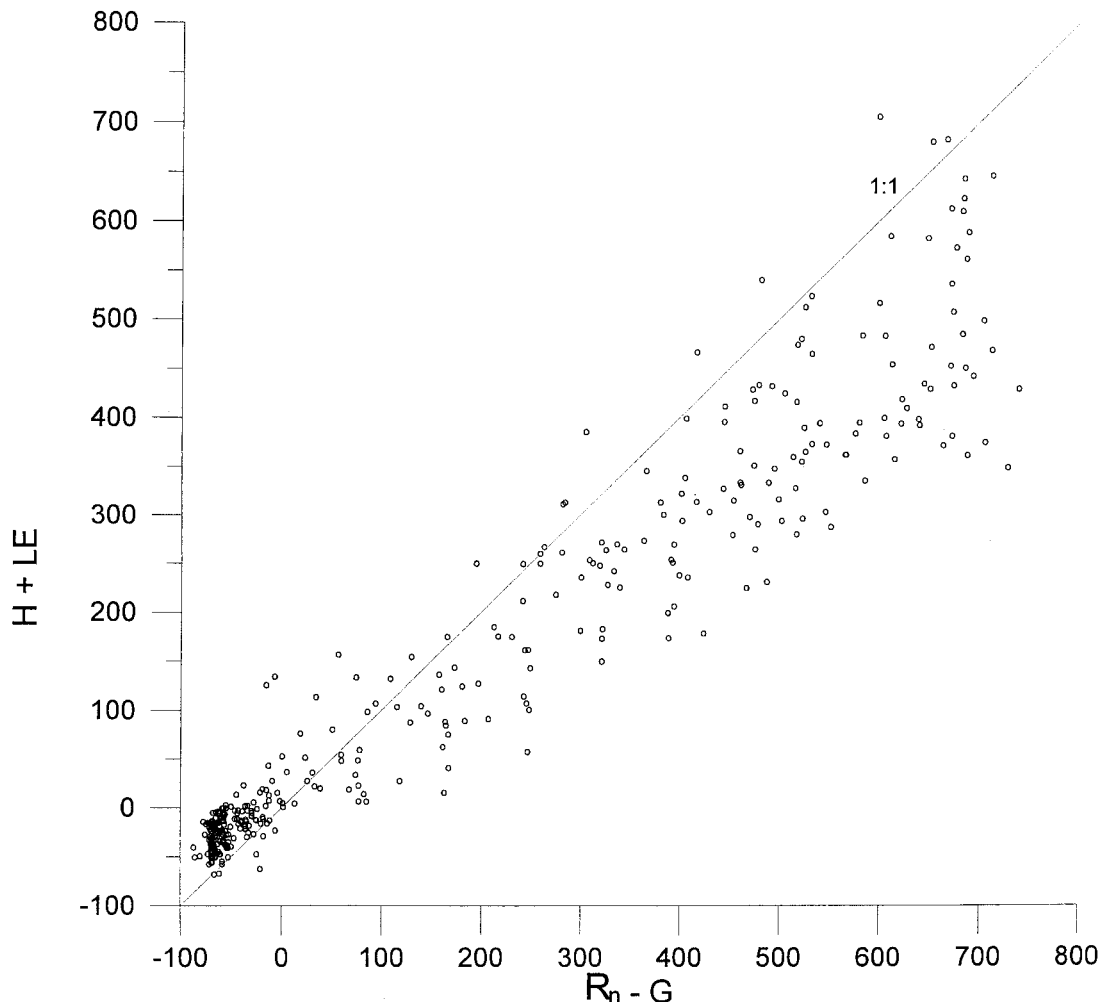


Fig. 7. Half-hour energy balance plot of  $H + LE$  versus  $R_n - G$  (open circles) for the period 30 July–8 August, 1996.

Fig. 8. This approach yielded an average nighttime (17:30–05:30) convection correction value of  $+0.08 \text{ mg m}^{-2} \text{ s}^{-1}$  and an average daytime (07:30 am to 13:00 pm) convection correction of  $-0.03 \text{ mg m}^{-2} \text{ s}^{-1}$  with a linear transition between the two values occurring from 05:30 to 07:30 MST and from 13:00 to 17:30 MST.

The corrected  $\text{CO}_2$  flux varies from  $-1.0 \text{ mg m}^{-2} \text{ s}^{-1}$  to near zero during daytime and averages about  $+0.25 \pm 0.2 \text{ mg m}^{-2} \text{ s}^{-1}$  at night. Note that daytime corrections are small. Thus, the solid line in Fig. 8 representing fully corrected  $\text{CO}_2$  flux is almost congruent with the dashed line of the LM correction during daytime. These flux results are similar to summertime  $\text{CO}_2$  exchange rates measured over a deciduous forest near Oak Ridge, TN ( $35^\circ\text{N}$ ,  $84^\circ\text{E}$ ) by Greco and Baldocchi (1996) who report daytime maximum of  $-1.1 \text{ mg m}^{-2} \text{ s}^{-1}$ , and night-time average of  $+0.18 \text{ mg m}^{-2} \text{ s}^{-1}$ . However, they are lower than daytime maxima of  $-0.35 \text{ mg m}^{-2} \text{ s}^{-1}$ , and night-time mean of  $+0.22 \text{ mg m}^{-2} \text{ s}^{-1}$  observed over a boreal black spruce forest in Saskatchewan, Canada ( $54^\circ\text{N}$ ,  $105^\circ\text{W}$ ) by Jarvis et al. (1997).

### 3.5.2. $\text{O}_3$ fluxes

The  $\text{O}_3$  flux record is presented at the bottom of Fig. 8 along with the  $\text{O}_3$  storage parameter.  $\text{O}_3$  deposition is shown as a downward flux. Corrections of the  $\text{O}_3$  flux for a closed-path system were small, i.e.  $\pm 0.015 (\mu\text{g m}^{-2} \text{ s}^{-1})$ . Maximum  $\text{O}_3$  uptake occurs midday and ranges from  $-0.5$  to  $-0.6 \mu\text{g m}^{-2} \text{ s}^{-1}$ . Deposition rates are relatively high during night-time hours of JD 215 and 216 when atmospheric mixing was weak but  $\text{O}_3$  concentration declined by more than 20 ppb, and again on the night of JD 217–218 when turbulent mixing was stronger than usual. The three brief, positive and counter-gradient fluxes of  $\text{O}_3$  (on JD 215, 216 and 217) occur at sundown during periods of atmospheric stability transition. These occurrences are associated with increase in relative humidity and vertical  $\text{O}_3$  concentration gradient, decrease in air temperature, wind speed and  $\text{O}_3$  concentration and (except on day 216) coincide with brief positive perturbations in both  $R_n$  and  $\int_0^h \frac{\partial \bar{c}}{\partial z} dz$ . Similar but weaker flux changes also occur on JD 213 and 214 at the same time of day. Summertime counter-gradient  $\text{O}_3$  fluxes have previously been reported at this site (Zeller and Hehn, 1994, 1996).

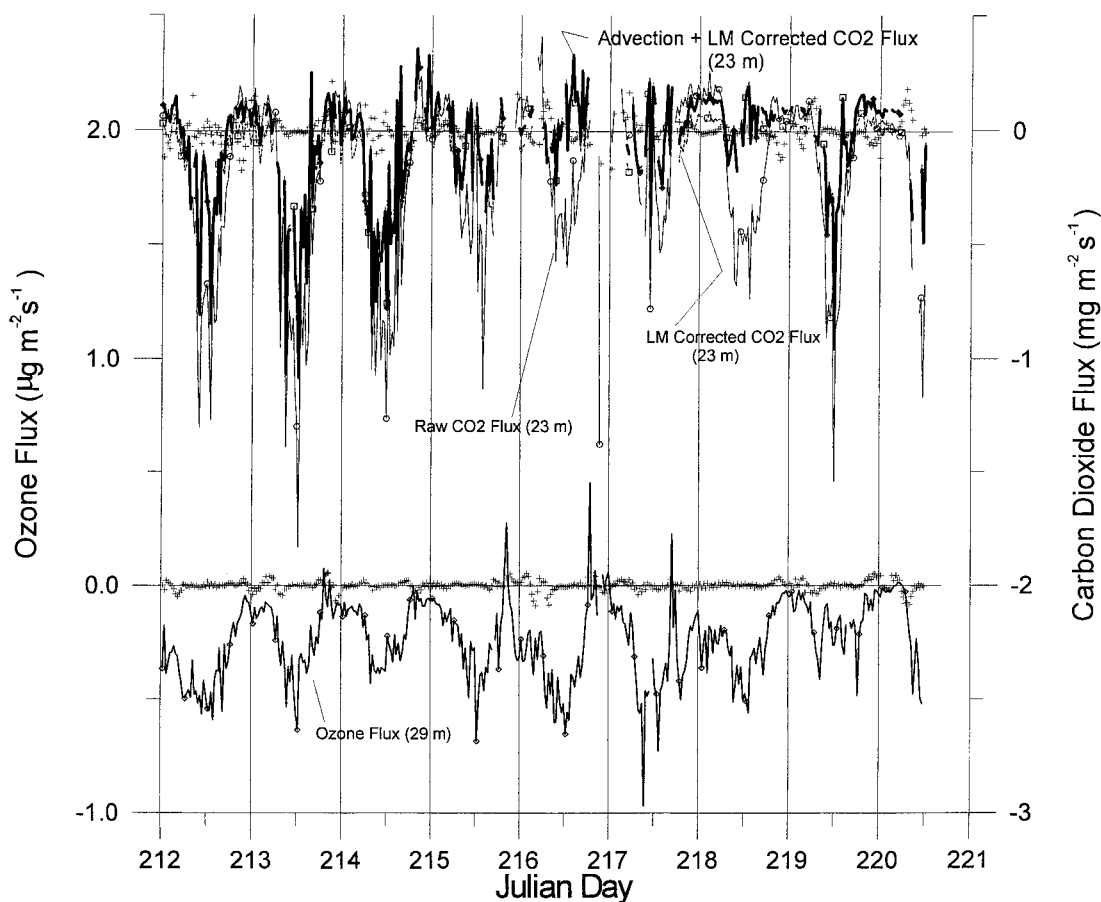


Fig. 8. Half-hour average 23 m  $\text{CO}_2$  flux: (1) as measured (raw) (thin solid line, open circle), (2) with sensible and latent heat LM (Leuning and Moncrieff, 1990) corrections (dashed line, open square), and (3) with LM correction and advection correction (heavy solid line, open diamond); 29 m  $\text{O}_3$  flux (solid line, open diamond); and local time-rate-of-change value for  $\text{CO}_2$  and  $\text{O}_3$  (plus signs) for the period 30 July–8 August, 1996.

### 3.6. Comparison of measured fluxes with model simulations

The FORFLUX model was run for the nine measurement days using observed standard weather data at GLEES, published site characteristics from Musselman (1994) (Table A2) and derived species physiological parameters (Table A1). Due to the lack of information, values of physiological parameter for Engelmann spruce were estimated from corresponding data for Douglas fir (Nikolov, 1997).

Figs. 9 and 10 illustrate the comparison of modeled and measured fluxes of  $\text{CO}_2$ ,  $\text{O}_3$  and latent heat. The model closely mimics daytime  $\text{CO}_2$  flux on days 212–214. However, it predicts higher values at night. On the night of JD 219–220, it closely mimics the  $\text{CO}_2$  flux while somewhat approximating the JD 219 daytime flux. It also does a good job for the night of JD 217–218 when a strong turbulent mixing prevailed. However, on days 215–218 and 220, measured and modeled daytime  $\text{CO}_2$  fluxes are quite different. Although we do not assume that the model is perfectly accurate, it is likely that the eddy covariance system is just not properly

registering the actual  $\text{CO}_2$  flux during these times. This conclusion is supported by the fact that a substantial downward flux of  $\text{O}_3$  is both modeled and measured on days 215–218. The magnitude of this  $\text{O}_3$  flux suggests a high canopy stomatal conductance and, hence, a high  $\text{CO}_2$  uptake during these days. In other words, the measured corrected  $\text{CO}_2$  flux does not seem to be consistent with the observed  $\text{O}_3$  deposition on days 215–218. The correlation between modeled and measured hourly  $\text{CO}_2$  fluxes for the entire sampling period is  $r^2=0.29$  with a slope of  $0.7\pm0.06$ ; intercept of  $-0.1\pm0.25$ ; and root mean square error (RMSE) of 0.27. The same statistics for days 212–214 is  $r^2=0.58$ , slope  $0.8\pm0.06$ , intercept  $0.01\pm0.2$  and RMSE 0.219. Since the  $\text{CO}_2$  measurements were made 6 m above the canopy, i.e. well within the turbulent mixing layer, these results corroborate the observations of Lee (1998) and Baldocchi (personal communication) for the need to measure  $\text{CO}_2$  fluxes close to the canopy top when in rough terrain (Lee et al., 1996). It also appears that it is necessary to sample  $\text{CO}_2$  concentrations and wind parameters near the forest floor and at several heights within the canopy to establish scalar profiles and in

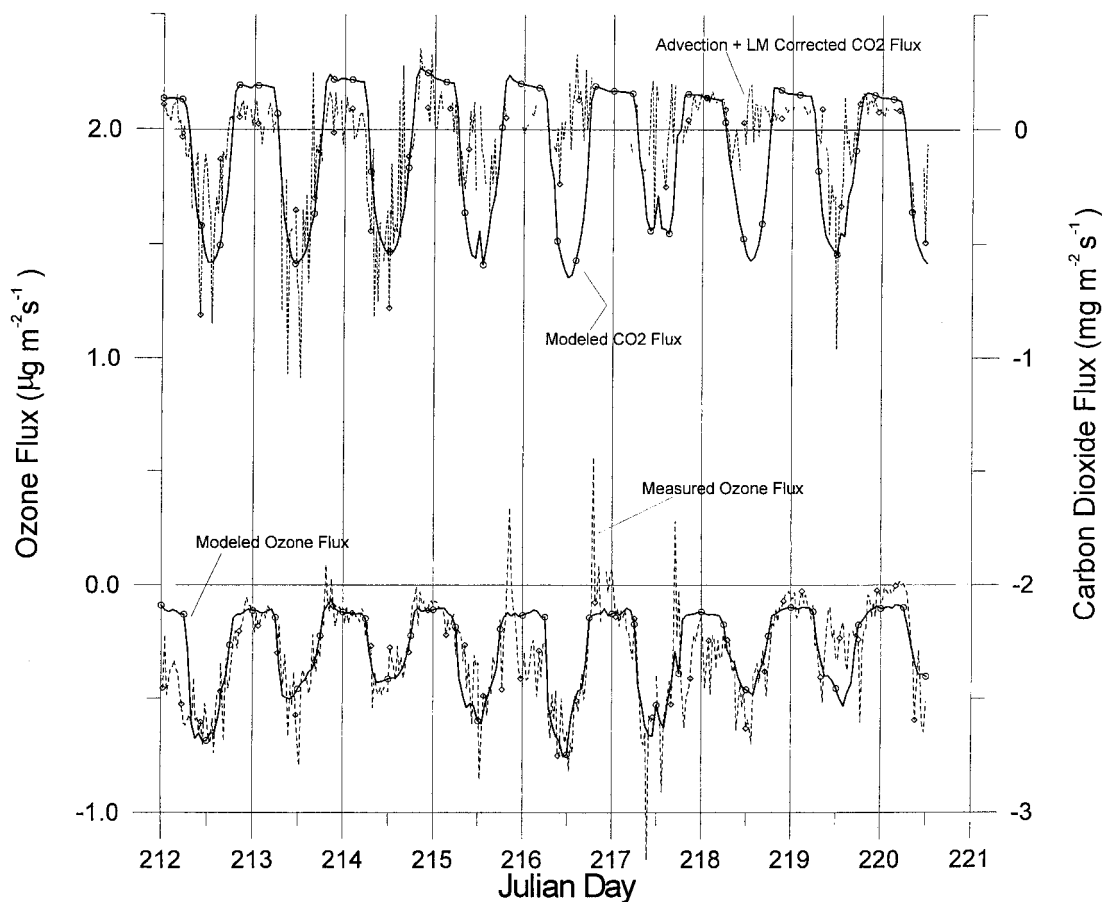


Fig. 9. Half-hour average 23 m  $\text{CO}_2$  flux: (1) with LM correction (Leuning and Moncrieff, 1990) and advection correction (dashed line, open diamond) and (2) modeled (solid line, open circle); and 29 m  $\text{O}_3$  flux (1) as measured with corrections (dashed line, open diamond), and (2) modeled (solid line, open circle) for the period 30 July–8 August, 1996.

order to adequately correct fluxes for advection and storage effects.

The comparison between modeled and measured  $O_3$  fluxes for the 9-day period is significantly better than that for  $CO_2$  ( $r^2=0.54$ ; slope,  $0.64 \pm 0.03$ ; intercept,  $-0.09 \pm 0.13$ ; RMSE, 0.151). The model successfully tracks observations during most hours. The average measured  $O_3$  flux for the 9-day period is  $0.30 \mu g m^{-2} s^{-1}$ , which compares rather well with  $0.28 \mu g m^{-2} s^{-1}$  predicted by the model. Apparently, sampling height is not as much a problem for  $O_3$  as it is for  $CO_2$ . Nevertheless, previously reported positive fluxes of  $O_3$  during daytime in the summer by Zeller and Hehn (1996) and the occurrence of occasional positive fluxes in this study, point to the need to correct for possible  $O_3$  advection as well.

Fig. 10 presents the comparison between modeled latent heat flux and measured LE at 12 m above the canopy. Observed daytime maximum values range between 300 and  $400 W m^{-2}$ . These fluxes are somewhat

higher compared to  $200\text{--}300 W m^{-2}$  measured over a deciduous forest by Greco and Baldocchi (1996) and  $150\text{--}200 W m^{-2}$  observed over a boreal black spruce forest by Jarvis et al. (1997). The statistics for the agreement between hourly modeled and measured LE fluxes are:  $r^2=0.69$ ; slope,  $0.78 \pm 0.03$ ; intercept,  $21.9 \pm 50$ ; RMSE, 53.8. The model predicted an average latent heat flux of  $99.6 W m^{-2}$  for the sampling period. This value is virtually identical with a mean LE of  $99.7 W m^{-2}$  yielded by the measurements. It appears that sampling height is also not an issue when measuring turbulent water vapor exchange at GLEES.

Based on the above results, we believe that the FOR-FLUX model adequately predicted ecosystem-atmosphere exchange of water vapor and  $O_3$  for the nine sampling days at the study site. Due to uncertainties in the  $CO_2$  flux measurements, it is not possible to directly assess the modeled carbon exchange. However, since canopy stomatal conductance is known to exert strong control on all the studied fluxes, and the model

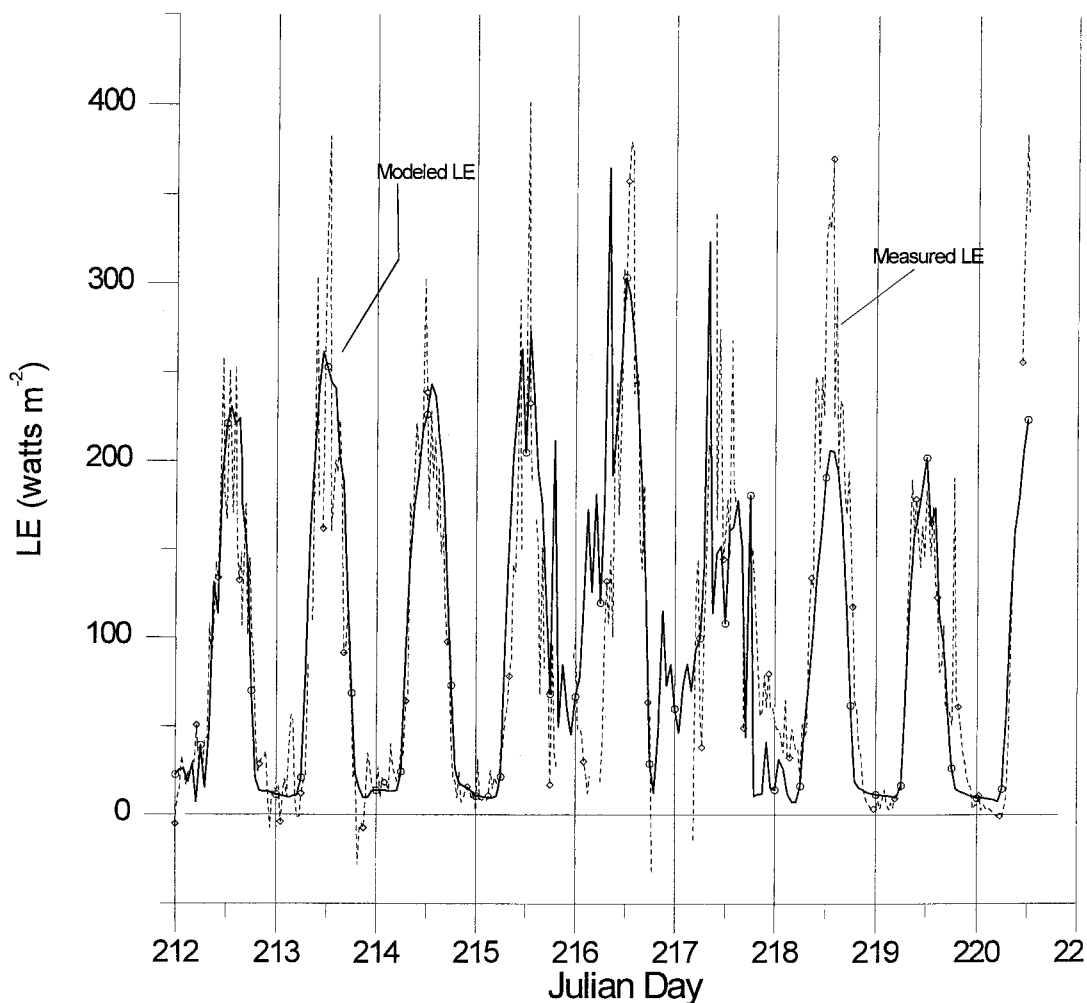


Fig. 10. Half-hour average 29 m latent heat (LE) flux: (1) as measured with corrections (dashed line, open diamond), and (2) modeled (solid line, open circle) for the period 30 July–8 August, 1996.

adequately reproduced evapotranspiration and total O<sub>3</sub> deposition, we assume that it also satisfactorily simulated the CO<sub>2</sub> flux.

### 3.7. Application of the FORFLUX model to assess annual fluxes at GLEES

The FORFLUX model was used to estimate seasonal dynamics and annual totals of major CO<sub>2</sub>, water vapor and O<sub>3</sub> fluxes over this pristine ecosystem. The model was run for a full year using standard hourly meteorological data (i.e. temperature, relative humidity, solar radiation, precipitation) measured at the Brooklyn Tower in 1996, and O<sub>3</sub> concentration recorded for the same time period at a nearby EPA NDDN site. The site-specific and species-physiological parameters required by the model were the same ones used for the verification procedure above (Tables A1 and A2). Based on observations, an initial snowpack depth of 1.6 m was assumed to be present on the site on 1 January 1996. The initial mean soil moisture content was set to 70% of the field capacity value. The latter is estimated in the model from soil texture information.

Fig. 11 depicts simulated seasonal dynamics of the total ecosystem–atmosphere exchange of CO<sub>2</sub>, water vapor and O<sub>3</sub> at GLEES for 1996. The model predicts a

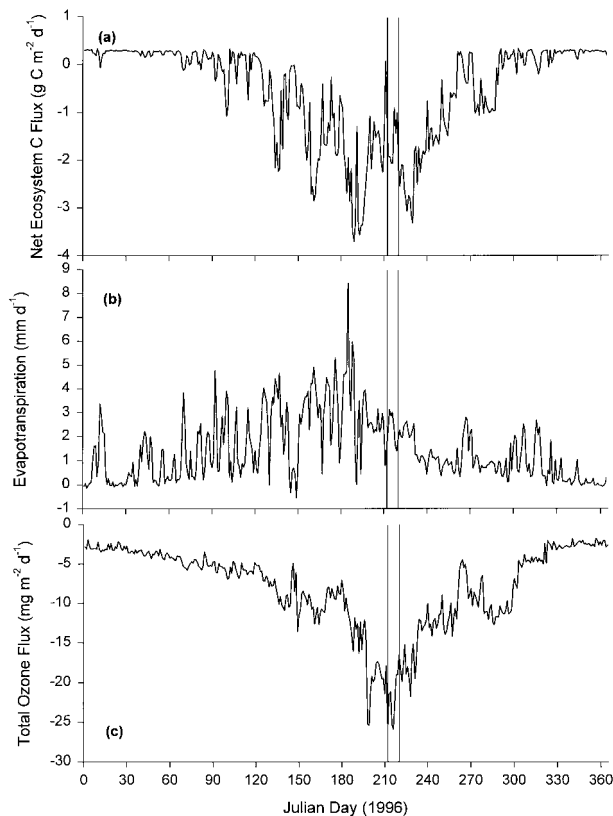


Fig. 11. Seasonal dynamics of daily net ecosystem–atmosphere exchange (NEE) at GLEES predicted by the FORFLUX model in 1996: (a) CO<sub>2</sub> uptake; (b) evapotranspiration; and (c) O<sub>3</sub> deposition.

large fluctuation of daily fluxes for all variables due to variations in meteorological conditions. The daily net ecosystem carbon exchange is estimated to vary from  $-3.71$  to  $+0.34$  g C m<sup>-2</sup> day<sup>-1</sup> over the growing season having an average value of about  $-0.53$  g C m<sup>-2</sup> day<sup>-1</sup> (Fig. 11a). A significant reduction in ecosystem carbon uptake is predicted around day 210 due to an increased CO<sub>2</sub> efflux from soil in response to elevated ground temperature and soil moisture following the snowmelt. Daily evapo-transpiration (ET) ranges from  $-0.55$  to  $8.42$  mm day<sup>-1</sup> with negative values caused by vapor condensation on the snow surface (Fig. 11b). ET is noticeable for a smaller seasonal amplitude compared to CO<sub>2</sub> flux. This is a result of a substantial snowpack sublimation during winter and a modest canopy transpiration in the summer.

The total O<sub>3</sub> deposition follows a seasonal pattern, which is distinctly different from that of carbon and water fluxes (Fig. 11c). Daily predicted O<sub>3</sub> uptake by the ecosystem gradually increases from January through to the end of June, and then suddenly leaps by about 100%, staying high throughout the end of August. The

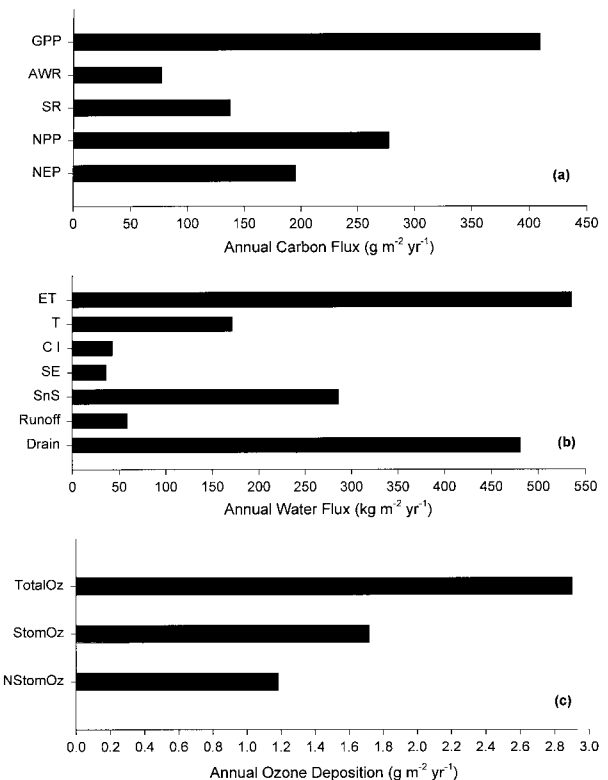


Fig. 12. Component annual fluxes of (a) carbon, (b) water, and (c) ozone predicted by the FORFLUX model at GLEES in 1996 (GPP, net canopy photosynthesis; AWR, Above-ground woody respiration; SR, soil respiration; NPP, net primary production; NEP, net ecosystem production; ET, evapo-transpiration; T, transpiration; CI, canopy rainfall interception; SE, soil evaporation; SnS, snow sublimation; runoff, surface runoff; drain, subsurface drainage; TotalOz, total ozone deposition; StomOz, ozone uptake by leaf stomata; NStomOz, ozone deposition to non-transpiring plant surfaces).

rate of predicted  $O_3$  deposition is relatively low in winter and spring due to the presence of a snowpack on the ground. Snow possesses about a 10-time higher resistance to  $O_3$  exchange than soils. Snowpack also prevents soil from warming up in early spring, which has a negative impact on plant stomatal conductance in the model reducing the  $O_3$  uptake by vegetation. The daily  $O_3$  deposition rate at GLEES varies from 1.8 to 25.9  $mg\ m^{-2}\ day^{-1}$  over the course of the year. The magnitude of  $O_3$  fluxes predicted in summertime is consistent with measurements made by Zeller and Hehn (1996) in July of 1992. The subalpine ecosystem was predicted to destroy  $O_3$  at an average rate of 7.92  $mg\ m^{-2}\ day^{-1}$  in 1996.

Fig. 12 presents component annual fluxes of the three scalars predicted by the FORFLUX model at GLEES in 1996. Net ecosystem–atmosphere exchange of carbon (NEP) is 195  $g\ m^{-2}\ year^{-1}$  while net primary production of vegetation (NPP) is 277  $g\ m^{-2}\ year^{-1}$ . NEP is less than NPP because it includes the heterotrophic soil respiration. These rather low productivity values are in agreement with the observed slow growth rate of trees in the Medicine Bow Mountains (Earle, 1993). Predicted NPP is also consistent with values reported for boreal forest stands (e.g. Moore, 1996; Ryan et al., 1997). The low NPP and retarded tree growth are attributed to the harsh alpine environment characterized by a cool and short growing season. NEP and NPP are estimated to be 47.8 and 67.7% of the net canopy photosynthesis (or GPP), respectively.

The modeled cumulative ET is 535  $mm\ year^{-1}$  (Fig. 12b), which represents about 68% of the annual precipitation. Transpiration (174  $mm\ year^{-1}$ ) is only 32% of ET, and amounts to about 60% of the annual sublimation from snow. Surface runoff totals 59  $mm\ year^{-1}$ , most of which occurs during a short period in early summer after a rapid snow melt. The model suggests that most of the water released to the atmosphere by this subalpine ecosystem in 1996 came from snow and soil (i.e. ground evaporation represents ca. 60% of ET). Due to a relatively shallow soil column, the runoff is heavily dominated by subsoil drainage.

The model predicted a total  $O_3$  deposition of 2.9  $g\ m^{-2}\ year^{-1}$  to GLEES in 1996. About 59% of this amount, or 1.72  $g\ O_3\ m^{-2}\ year^{-1}$  is projected to be actively taken up by the foliage via leaf stomata. This is the  $O_3$  that comes into direct contact with leaf mesophyll cells and can adversely affect plant photosynthesis and growth. A flux of 1.18  $g\ O_3\ m^{-2}\ year^{-1}$  is estimated to be deposited to non-transpiring surfaces such as tree stems, branches, snow and soil.

The relatively large portion of non-stomatal  $O_3$  uptake (i.e. about 41% of the total flux) suggests that, when assessing pollution-induced vegetation damage, one needs to relate plant response to the *actively assimilated*  $O_3$  by foliage rather than to total deposition. In this regard, we propose a practical index for quanti-

fying vegetation vulnerability to  $O_3$  injury called Physiological  $O_3$  Uptake Per Unit of Leaf Area (POUPULA). It is a ratio of the time-integrated canopy  $O_3$  uptake by plant stomata to the one-sided LAI of the stand. Musselman and Massman (1999) proposed a similar flux-based index to assess  $O_3$  effects on vegetation. The operational implementation of their index, however, is currently impaired by the lack of quantitative information about plant defense mechanisms. POUPULA is not directly measurable on the field but it can be quantified using process models or derived from observed total  $O_3$  deposition amounts. Thus, we estimate POUPULA to be 0.614  $g\ O_3\ m^{-2}\ leaf\ area\ year^{-1}$  for GLEES in 1996. Since this index depends on both species physiology and the physical environment, its value is expected to vary among sites and between years on a particular site. It is, therefore, a future research goal to determine the spatial and temporal variability of this index and to interpret its values in terms of critical thresholds of  $O_3$  uptake for important vegetation types.

#### 4. Conclusion

Measuring carbon exchange has proven to be problematic at GLEES. This is partly because the eddy covariance technique employed above the canopy tends to undermeasure the nighttime ecosystem respiration. Three possible mechanisms have been proposed to explain the poor detection of respiratory  $CO_2$  flux above the canopy: (1) cool, heavier air containing  $CO_2$  may sink to the forest floor at night and advect/meander away under stable atmospheric conditions (Baldocchi et al., 1996; Lee et al., 1996); (2) vertical waves of air associated with stable wind shear can cause advection of pockets of  $CO_2$  (Lee et al., 1996); and (3) the possible inability of the data correction algorithms to accurately account for the 3-dimensional  $CO_2$  flow in mountainous terrain.

The corrected  $CO_2$  flux measured at the GLEES Brooklyn tower shows a reasonable late-summer behavior on certain days and some nights but rarely for a complete 24-h period. The maximum daytime  $CO_2$  fluxes over this subalpine spruce–fir forest ecosystem are  $-0.7 \pm 0.1\ mg\ m^{-2}\ s^{-1}$  while night-time exchange varied between 0.2 and  $0.3 \pm 0.2\ mg\ m^{-2}\ s^{-1}$ . This experiment has demonstrated the necessity to address  $CO_2$  advection and storage in future measurements at this site. An empirical correction adapted from Lee (1998) (Baldocchi, personal communication) is employed here in an attempt to address the night-time flux measurement problem. We have found that biophysical models might be useful when assessing eddy flux measurements made in rough terrain. Future  $CO_2$  flux measurements at GLEES should be made closer to the canopy top. Such measurements will also require a detailed study of the diel  $CO_2$  profile changes below the sampling height to



determine how many additional CO<sub>2</sub> concentration measurements must be made to accompany the flux corrections for advection and storage and to provide a valid data set for model validation. Our experience at GLEES suggests that the goal of accurately measuring net CO<sub>2</sub> flux should be either limited to very carefully chosen sites, or require an extensive set of simultaneously collected auxiliary CO<sub>2</sub> and meteorological data.

Both model simulations and measurements suggest that, in early August of 1996, the maximum daytime O<sub>3</sub> deposition rates were  $-0.5$  to  $-0.6 \mu\text{g m}^{-2} \text{s}^{-1}$ . These results compare well with peak fluxes of  $-0.3$  to  $-0.6 \mu\text{g m}^{-2} \text{s}^{-1}$  reported for GLEES in July 1992 by Zeller and Hehn (1996). The choice of above canopy measurement height does not appear to be critical for O<sub>3</sub> fluxes. However, previous O<sub>3</sub> flux measurements indicate that advection should also be studied, particularly if the occasionally measured counter-gradient fluxes are to be understood.

The measured sensible and latent heat fluxes are reasonable and agree well with model predictions. They provide some measure of confidence given the obtained energy balance closure. Future experiments at this site should account for heat storage within the canopy.

The FORFLUX model predicted the following annual fluxes for GLEES in 1996:  $195 \text{ g m}^{-2} \text{ year}^{-1}$  NEP;  $277 \text{ g m}^{-2} \text{ year}^{-1}$  NPP;  $535 \text{ mm year}^{-1}$  ET; and  $2.9 \text{ g m}^{-2} \text{ year}^{-1}$  total O<sub>3</sub> deposition (with an estimated uptake by leaf stomata of  $1.72 \text{ g O}_3 \text{ m}^{-2} \text{ year}^{-1}$ ). The potential ability of the model to discriminate between

total O<sub>3</sub> deposition and physiological O<sub>3</sub> uptake by foliage could provide a valuable tool for researchers who seek to relate pollution-induced vegetation damage to specific uptake amounts rather than to ambient O<sub>3</sub> concentration. In this regard, we introduce for more testing POUPULA, as a practical index for assessing vegetation vulnerability to O<sub>3</sub> injury. This indicator quantifies the amount of O<sub>3</sub> that comes into direct contact with plant photosynthetic machinery and, thus, is mostly responsible for the observed decline in vegetation productivity and crop yield.

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## Appendix A

### FORFLUX model input parameters and values used

Table A1

Input species physiological parameters used in the FORFLUX model for Engelmann spruce (*Picea engelmannii*)<sup>a</sup>

Parameter	Symbol	Units	Value
Maximum carboxylation velocity at 25°C	$V_{m25}$	$\text{mol m}^{-2} \text{s}^{-1}$	59.0
Light-saturated potential rate of electron transport	$J_{m25}$	$\text{mol m}^{-2} \text{s}^{-1}$	108.0 at 25°C
Activation energy for electron transport	$E_j$	$\text{J mol}^{-1}$	37,000
Kinetic parameter for CO <sub>2</sub> at 25°C	$K_{c25}$	$\text{mol mol}^{-1}$	$22 \times 10^{-5}$
Kinetic parameter for O <sub>2</sub> at 25°C	$K_{o25}$	$\text{mol mol}^{-1}$	0.41
Photosynthetic light loss factor	$f$	Decimal fraction	0.50
Composite stomatal sensitivity	$m$	Dimensionless	14.20
Empirical constant for stomatal conductance	$b_{sv}$	$\text{mol m}^{-2} \text{s}^{-1}$	0.019
Mean needle width	$d$	m	$1.1 \times 10^{-3}$
Individual leaf reflectance in the visible wavelength	$\alpha_{f, \text{vis}}$	Decimal fraction	0.09
Individual leaf transmittance in the visible wavelength	$\tau_{f, \text{vis}}$	Decimal fraction	0.04
Individual leaf reflectance in the near infrared wavelength	$\alpha_{f, \text{nir}}$	Decimal fraction	0.50
Individual leaf transmittance in the near infrared wavelength	$\tau_{f, \text{nir}}$	Decimal fraction	0.39
Root conductance parameter	$a_r$	Dimensionless	2.0
Root conductance parameter	$b_r$	Dimensionless	30.0
Root conductance parameter	$c_r$	Decimal fraction	0.25
Critical temperature for root conductance	$T_{cr}$	°C	4.5
Root signal parameter	$A_r$	$\mu\text{mol kg m}^{-6} \text{h}^{-1}$	$12 \times 10^{-6}$
Root signal parameter	$B_r$	$\text{m}^{-1}$	0.026
Root signal parameter	$C_r$	$\text{kg m}^{-2} \text{h}^{-1}$	0.014

<sup>a</sup> For parameter explanation, see Nikolov (1997).

Table A2

Site-specific input parameters of the FORFLUX model for GLEES

Parameter	Symbol	Units	Value
Slope inclination	$I_s$	Degree	0°
Latitude	$L_g$	Deg. Min	41°22'N
Altitude	$Z_a$	m asl	3186
Slope azimuth	$A_z$	Degree	0°
Ambient CO <sub>2</sub> concentration	$C_a$	μmol mol <sup>-1</sup>	363
Ambient O <sub>2</sub> concentration	$O_a$	mol mol <sup>-1</sup>	0.29
Mean annual air temperature	$T_{m,o}$	°C	3.0
Mean annual temperature amplitude	—	°C	10.0
Total canopy leaf area index	$L_{tot}$	m <sup>2</sup> m <sup>-2</sup>	2.8
Canopy foliage clumping factor	$\Omega_{cl}$	Decimal fraction	0.70
Mean leaf inclination angle	$\Theta$	Degree	57.4
Above-ground sapwood N	$N_s$	g m <sup>-2</sup>	3.20
Above-ground C allocation coefficient	$\delta_a$	Decimal fraction	0.50
Fraction of soil CO <sub>2</sub> efflux in root respiration	—	Decimal fraction	0.40
Total soil depth	$Z_{max}$	m	0.70
Rooting depth	$Z_r$	m	0.52
Soil bulk density	$\rho_{bs}$	mg m <sup>-3</sup>	1.20
Soil sand content	Sa	% of dry weight	47
Soil clay content	Cl	% of dry weight	13
Soil organic C content	C	% of dry weight	0.5
Soil rock content	$f_r$	% of volume	50
Visible soil reflectance	$\alpha_{s, vis}$	Decimal fraction	0.24
Near-infrared soil reflectance	$\alpha_{s, vis}$	Decimal fraction	0.30
Soil roughness length	$z_o$	m	0.01
Water table presence	—	Boolean	False
Initial snow depth	—	m	0.0
Initial surface soil moisture	—	% of field cap.	65
Initial bottom soil moisture	—	% of field cap.	85

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